

Localization of the ground state of the Nelson model

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- FKF for the renormalized Nelson model
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Let us consider the Schrödinger equation:

$$hf = Ef, \quad h = -\frac{1}{2}\Delta + V.$$

Spatial decays of f

(1) L^2 -estimate: $\|e^{c|x|}f\| < \infty$

(2) point-wise-estimate: $a_1e^{-c|x|} \leq |f(x)| \leq a_2e^{-c|x|}$ a.e. $x \in \mathbb{R}^3$.

We are interested in "point-wise-estimate" from both upper and lower.
Roughly speaking solving the equation ($d = 1$)

$$-\frac{1}{2}(e^{-r})'' + Ve^{-r} = Ee^{-r}$$

we obtain $(r')^2 = 2(V - E)$, $r(x) = \int^x \sqrt{2(V(x) - E)}dx \sim |x|\sqrt{2(V(x) - E)}$.

$$|f(x)| \sim Ae^{-c|x|\sqrt{V(x)}}.$$

FKF

$$(e^{-th} f)(x) = \mathbb{E}^x[e^{-\int_0^t V(B_s) ds} f(B_t)], \text{ a.e. } x \in \mathbb{R}^3$$

► $(B_t)_{t \geq 0}$ is 3-dim Brownian motion with $B_0 = x$ on $(\Omega, \mathcal{Y}, W^x)$.

$$\mathbb{E}^x[\dots] = \int_{\Omega} \dots dW^x.$$

Ground state $h\varphi_p = E\varphi_p$, $E = \inf \text{Spec}(h)$.

$$\frac{e^{-th} f}{\|e^{-th} f\|} \rightarrow \varphi_p \quad (t \rightarrow \infty) \text{ in } L^2, \quad (f \geq 0)$$

► We keep in mind that

$$\varphi_p(x) \sim \frac{(e^{-th} f)(x)}{\|e^{-th} f\|} = \frac{\mathbb{E}^x[e^{-\int_0^t V(B_s) ds} f(B_t)]}{(\int |\mathbb{E}^x[e^{-\int_0^t V(B_s) ds} f(B_t)]|^2 dx)^{1/2}} \quad t \gg 1$$

Examples $h = -\frac{1}{2}\Delta + \frac{1}{2}|x|^2$

►(1) Set $hf = Ef$. Then $|x|\sqrt{V(x)} = |x|^2$ and it is known that

$$f(x) = h(x)e^{-|x|^2/2}.$$

►(2) Set $f = \mathbb{1} \notin L^2$. Then

$$\begin{aligned} \frac{(e^{-th}\mathbb{1})(x)}{\|e^{-th}\mathbb{1}\|} &= \frac{\mathbb{E}^x[e^{-\frac{1}{2}\int_0^t |B_s|^2 ds}]}{(\int |\mathbb{E}^x[e^{-\frac{1}{2}\int_0^t |B_s|^2 ds}]^2 dx)^{1/2}} \\ &= \frac{e^{-|x|^2/2\coth t}}{(\int e^{-|x|^2/\coth t} dx)^{1/2}} \quad (\text{Yor formula}) \\ &\xrightarrow{t \rightarrow \infty} \pi^{-3/4} e^{-|x|^2/2} \quad (\text{ground state}) \end{aligned}$$

Let $V^-(x) = \inf\{V(y) \mid |x - y| \leq 1\}$ and $hf = Ef$.

$$f(x) = e^{-t(h-E)} f(x) = \mathbb{E}^x[e^{-\int_0^t (V(B_s) - E) ds} f(B_t)]$$

Upper bound (Carmona 79)

$$|f(x)| \leq ae^{-c|x|}\sqrt{V_-(x)-E}$$

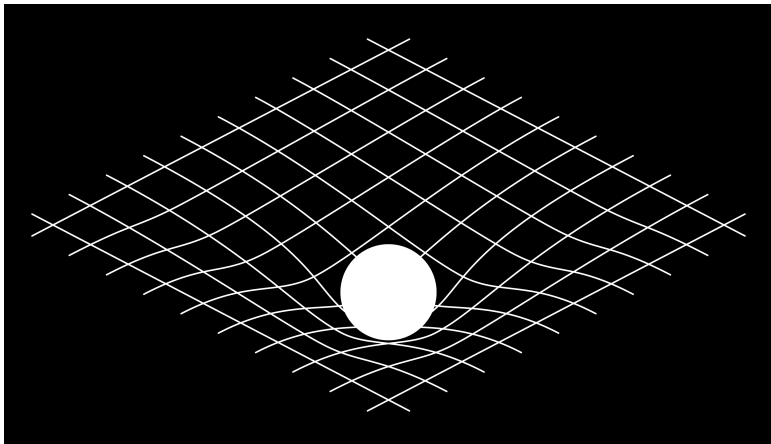
Let φ_p be the ground state. Let $K = [-b_1, b_1] \times [-b_2, b_2] \times [-b_3, b_3]$.

$$\inf_{x \in K} \varphi_p(x) > \exists \varepsilon_K > 0.$$

※ If $V = V_+ - V_-$ and V_+ is local-Kato, then $x \mapsto \varphi_p(x)$ is cont due to the smoothing effect of $e^{t\Delta}$. $V_a(x) = \sup\{V(y) \mid |x_j - y_j| \leq 1 + |x_j|, j = 1, 2, 3\}$

Lower bound (Carmona 79)

$$|\varphi_p(x)| \geq a\varepsilon_K e^{-c|x|}\sqrt{V_a(x)-E}.$$



$$H_{\varepsilon, \Lambda} = H_S \otimes \mathbb{1} + \mathbb{1} \otimes H_f + g\phi(\rho_\Lambda(\cdot - x)) - E_\Lambda$$

Gaussian random variable $\phi(f), (f \in L^2_{\mathbb{R}}(\mathbb{R}^3))$ on (Q, Σ, μ)

$$\mathbb{E}_{\mu}[\phi(f)] = 0, \quad \mathbb{E}_{\mu}[\phi(f)\phi(g)] = \frac{1}{2}(f, g)_2$$

$$\ast \mathbb{E}_{\mu}[\dots] = \int_Q [\dots] d\mu \text{ and } \mathbb{E}_{\mu}[e^{z\phi(f)}] = e^{\frac{z^2}{4}\|f\|_2^2}$$

Boson Fock space

$$\mathcal{F} = L^2(Q) = \overline{LH\{\phi(f_1) \cdots \phi(f_n)\}}$$

▶ $H_f = d\Gamma(\hat{\omega})$ is **"the free field Hamiltonian"** defined by

$$H_f : \phi(f_1) \cdots \phi(f_n) := \sum_j : \phi(f_1) \cdots \phi(\hat{\omega}f_j) \cdots \phi(f_n) :$$

$$H_f \mathbb{1} = 0$$

where $\hat{\omega} = \omega(-i\nabla)$ with

$$\omega(k) = \begin{cases} |k| & \text{massless} \\ \sqrt{|k|^2 + \mu^2} & \text{massive} \end{cases}$$

Total Hilbert space

$$\mathcal{H} = L^2(\mathbb{R}^3) \otimes \mathcal{F} \cong L^2(\mathbb{R}_x^3 \times Q_\phi)$$

$$\ast F = F(x, \phi) \in \mathcal{H}$$

Nelson Hamiltonian with UV cutoff Λ and IR cutoff ε

$$H_{\varepsilon, \Lambda} = H_S \otimes \mathbb{1} + \mathbb{1} \otimes H_f + g\phi(\rho_\Lambda(\cdot - x)) - E_\Lambda$$

$$\blacktriangleright H_S = -\frac{1}{2}\Delta_x + V$$

$\blacktriangleright g \in \mathbb{R}$ coupling constant

\blacktriangleright **Interaction:** $\phi(\rho_\Lambda(\cdot - x))$ Gaussian r.v. with test function $\rho_\Lambda(\cdot - x)$

► ρ_Λ has two parameters $0 \leq \varepsilon < \Lambda < \infty$. IR cutoff ε . UV cutoff Λ .
The Fourier transform of $\rho_\Lambda(\cdot - x)$ is given by

$$\hat{\rho}_\Lambda(k, x) = \frac{e^{-ikx}}{\sqrt{\omega(k)}} \mathbb{1}_{\varepsilon \leq |k| \leq \Lambda} \quad 0 \leq \varepsilon < \Lambda < \infty$$

► Removal of UV cutoff:

$$\lim_{\Lambda \rightarrow \infty} \hat{\rho}_\Lambda(k, x) = \hat{\rho}_\infty(k, x) = \frac{e^{-ikx}}{\sqrt{\omega(k)}} \mathbb{1}_{\varepsilon \leq |k|} \notin L^2(\mathbb{R}^3)$$

Then $\phi(\rho_\infty(\cdot - x))$ is not well-defined as a Gaussian r.v.

$$\text{Let } E_\Lambda = -\frac{1}{2} \int_{\mathbb{R}^3} \frac{\mathbb{1}_{\varepsilon \leq |k| \leq \Lambda}}{\omega(k)(\omega(k) + |k|^2/2)} dk (\rightarrow -\infty).$$

Thm (E. Nelson 64)

$$\exists \text{ s.a. } H_\infty > -\infty \text{ st } u\text{-}\lim_{\Lambda \rightarrow \infty} e^{-t(H_\Lambda - E_\Lambda)} = e^{-tH_\infty}.$$

► The **methods** of renormalization of the Nelson Hamiltonian

- Operator theory (Nelson 64)
- Stochastic method (Gubinelli-FH-Lőrinczi14, Matte-Møller17)

FKF for $\Lambda = \infty$

Gubinelli-FH-Lőrinczi (JFA14), Matte-Møller(PTRF17)

Let $F, G \in \mathcal{H}$. Then

$$(F, e^{-tH_\infty} G) = \int_{\mathbb{R}^3} \mathbb{E}^x \left[(F(B_0), e^{-\int_0^t V(B_s) ds} K_t G(B_t))_{\mathcal{F}} \right] dx.$$

Here the integral kernel is given by

$$K_t = e^S e^{a^\dagger(U_t)} e^{-tH_f} e^{a(\bar{U}_t)} \quad (a.s. \text{ bounded})$$

By FKF

$$\varphi_g(x) = \mathbb{E}^x \left[e^{-\int_0^t (V(B_s) - E) ds} K_t \varphi_g(B_t) \right] \quad a.e. (x, \phi) \in \mathbb{R}^3 \times Q.$$

φ_g : ground state $\iff H\varphi_g = E\varphi_g$ with $E = \inf \sigma(H)$

	$\varepsilon > 0$	$\varepsilon = 0$
$\mu > 0$	<i>exist</i>	<i>exist</i>
$\mu = 0$	<i>exist</i>	<i>not exist</i>

Figure: ε IR cutoff, μ mass, Ground state for $\Lambda \leq \infty$

▶ ($\Lambda < \infty$) Bach-Fröhlich-Sigal (AdvMath95), Arai-Hirokawa (JFA97), Spohn (LMP99), Gérard (AHP00), Griesemer-Lieb-Loss (InvMath01)

▶ ($\Lambda = \infty$) $|g| \ll 1 \rightarrow$ Hirokawa-FH-Spohn (Adv Math 05)

$\forall g \rightarrow$ FH-Matte, preprint (19)

Upper bound (FH20)

Suppose that $V = W - U$ st $U \geq 0$ and $U \in L^p$ for $3/2 < p < \infty$, $W \in L^1_{loc}$ and $\inf_x W > -\infty$. Then

$$\|\varphi_g(x)\|_{\mathcal{F}} \leq Ae^{-c|x|}\sqrt{W_-(x)}.$$

► Idea $\|\varphi_g(x)\|_{\mathcal{F}} \geq (\mathbf{1}, \varphi_g(x))_{\mathcal{F}} \stackrel{a.e.x}{=} \mathbb{E}^x[e^{-\int_0^t (V(B_s) - E) ds} (\mathbf{1}, K_t \varphi_g(x))_{\mathcal{F}}]$

Lower bound (FH20)

Let $V = W - U$ st $U \geq 0$, $U \in L^p$ for $3/2 < p < \infty$, and W is **local Kato**. Then the map $\mathbb{R}^3 \ni x \mapsto (\mathbf{1}, \varphi_g(x))$ is continuous, and

$$\|\varphi_g(x)\|_{\mathcal{F}} \geq Ae^{-c|x|}\sqrt{W_a(x)}.$$

► Let $0 \leq u \in L^2(\mathbb{R}^3)$ and since $(u \otimes \mathbb{1}, \varphi_g) \neq 0$,

$$u_t = \frac{e^{-tH_\infty} u \otimes \mathbb{1}}{\|e^{-tH_\infty} u \otimes \mathbb{1}\|_{\mathcal{H}}} \xrightarrow{t \rightarrow \infty} \varphi_g \text{ strongly in } \mathcal{H}$$

For $A : \mathcal{H} \rightarrow \mathcal{H}$, by FKF \exists a prob. measure μ_t on a path space st

$$(u_t, Au_t) = \mathbb{E}_{\mu_t}[f_A^t].$$

Hence

$$(\varphi_g, A\varphi_g) = \lim_{t \rightarrow \infty} \mathbb{E}_{\mu_t}[f_A^t].$$

Gibbs measure. FH AdvMath (14), FH-Matte (19)

$\exists \mu_\infty$ st $\mu_t \rightarrow \mu_\infty$ ($t \rightarrow \infty$) in the local sense.

Let N be the number operator.

Super exp-decay of bosons FH-Matte(19)

$$(\varphi_g, e^{+\beta N} \varphi_g) = \mathbb{E}_{\mu_\infty} \left[e^{+(1-e^{+\beta}) \int_0^\infty ds \int_{-\infty}^0 dr W(B_s - B_r, s-r)} \right] < \infty \quad \forall \beta \in \mathbb{R}.$$

Gaussian domination FH+Matte (19)

$$(\varphi_g, e^{+\beta \phi(f)^2} \varphi_g) = \frac{1}{\sqrt{1-\beta \|f\|^2}} \mathbb{E}_{\mu_\infty} \left[e^{+\frac{\beta (\int_{-\infty}^\infty S(B_s, s) ds)^2}{(1-\beta \|f\|^2)}} \right] < \infty \quad \forall \beta < 1/\|f\|^2$$

and

$$\lim_{\beta \uparrow 1/\|f\|^2} (\varphi_g, e^{+\beta \phi(f)^2} \varphi_g) = \infty$$

Cf. $h = -\frac{1}{2}\Delta + \frac{1}{2}|x|^2$ and $hf_n = Ef_n$. $\lim_{\beta \uparrow 1} (f_n, e^{\beta |x|^2} f_n) = \infty$.

Concluding remarks $\Lambda = \infty$

0. All results we obtain are **independent of coupling constants**.
1. Suppose that $\mu = 0$ (massless).
 - Let $\varepsilon > 0$. Then H_∞ has the ground state.
 - Let $\varepsilon = 0$. Then H_∞ has no ground state.
2. Properties of ground state:
 - **spatial decay**: $e^{-|x|\sqrt{V(x)}} \leq \|\varphi_g(x)\| \leq e^{-|x|\sqrt{V(x)}}$.
 - **super-exp. decay**: $\|e^{\beta N} \varphi_g\| < \infty$ for **all** $\beta \in \mathbb{R}$.
 - **Gaussian domination**: $\lim_{\beta \uparrow 1/\|f\|^2} (\varphi_g, e^{\beta \phi(f)^2} \varphi_g) = \infty$
3.
 - For decaying potential V we can also see a spatial decay.
 - Agmon metric