

# On the Embedded Eigenvalue of Relativistic Schrödinger Operator \*

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J. von-Neumann and E. Wigner showed that Schrödinger operator

$$H_{\text{NR}} = H_{\text{NR}}(V) := -\frac{1}{2}\Delta + V(x), \quad x \in \mathbb{R}^3 \quad (1)$$

may have a positive eigenvalue even if  $V(x)$  is a smooth decaying potential[2]. Actually, they constructed the following potential and eigenfunction

$$V_{\text{NW}}(x) := -32 \frac{\sin |x| (g(|x|)^3 \cos |x| - 3g(|x|)^2 \sin^3 |x| + g(|x|) \cos |x| + \sin^3 |x|)}{(1 + g(|x|)^2)^2}, \quad (2)$$

$$u_{\text{NW}}(x) := \frac{\sin |x|}{|x|(1 + g(|x|)^2)}. \quad (3)$$

where  $g(|x|) = 2|x| - \sin 2|x|$ . Then  $H_{\text{NR}}(V_{\text{NW}})u_{\text{NW}} = u_{\text{NW}}$  holds, and  $1 \in \sigma_{\text{p}}(H_{\text{NR}})$ . Since,  $V_{\text{NW}}$  is a decaying potential, the essential spectrum of  $H_{\text{NR}}$  is  $[0, \infty)$ . Hence, eigenvalue one is embedded in the essential spectrum.

In this talk, we generalize the Neumann and Wigner's result to the relativistic Schrödinger operator

$$H = (-\Delta + m^2)^{1/2} - m + V, \quad (4)$$

acting on  $L^2(\mathbb{R})$ , where  $m \geq 0$  is a mass of the particle. The operator  $(-\Delta + m^2)^{1/2}$  is defined through a functional calculus. As in the non-relativistic case, if  $V$  is decaying, then essential spectrum of  $H$  is  $[0, \infty)$ .

We define

$$h(x) := \frac{1}{1 + g(x)^2} \quad (5)$$

$$f(x) := \left( \sqrt{(-i\frac{d}{dx} + 1)^2 + m^2} + \sqrt{(-i\frac{d}{dx} - 1)^2 + m^2} \right) h(x). \quad (6)$$

and set

$$u(x) := f(x) \sin x \quad (7)$$

$$V(x) := \lambda - \frac{1}{u(x)} \left( \sqrt{-\frac{d^2}{dx^2} + m^2} - m \right) u(x), \quad (8)$$

where  $\lambda := \sqrt{1 + m^2} - m > 0$ . The following holds.

**Theorem 1.** *Let  $H$  be a relativistic Schrödinger operator with  $V$  defined by (8). If  $m \geq 146$ ,*

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then  $V$  is smooth potential with the property that  $V(x) = O(1/|x|)$  as  $|x| \rightarrow \infty$ . Moreover  $u \in D(H)$ , and  $\lambda$  and  $u$  satisfy the eigenvalue equation

$$Hu = \lambda u, \quad u \in D(H). \quad (9)$$

The above result can be extended to the three dimensional case:

**Theorem 2.** Assume  $m \geq 146$ . Let  $W(x) = V(|x|)$ ,  $x \in \mathbb{R}^3$  and define

$$\tilde{H} = \sqrt{-\Delta + m^2} - m + W(x) \quad (10)$$

which acts on  $L^2(\mathbb{R}^3)$ . Then

$$v(x) = \frac{u(|x|)}{\sqrt{4\pi|x|}}$$

belongs to  $D(\tilde{H})$ , and the eigenvalue equation  $\tilde{H}v = \lambda v$  holds.

**Remark** By restoring the speed of light  $c$  as a parameter in the operator, the eigenfunction and potential constructed above converge to the expressions obtained by von Neumann and Wigner as  $c \rightarrow \infty$ .

The above theorem doesn't cover the massless case  $m = 0$ . Instead of the construction of the positive eigenvalue, we can discuss the existence of the zero energy eigenvalue.

**Theorem 3.** Let us define the potential  $\tilde{V}_\nu$  and function  $v_\nu(x)$ .

$$\tilde{V}_\nu(x) := \begin{cases} -\frac{2(1-2\nu)\Gamma(\nu-\frac{1}{2})}{(1-\nu)\sqrt{\pi}\Gamma(\nu-1)}(1+x^2)^\nu {}_2F_1\left(2, \frac{1}{2} + \nu; \frac{3}{2}; -x^2\right), & \text{if } \nu \neq 1 \\ -\frac{2}{1+x^2}, & \text{if } \nu = 1, \end{cases}$$

$$v_\nu(x) := \frac{x}{(1+x^2)^\nu}.$$

Then the eigenvalue equation  $\sqrt{-\frac{d^2}{dx^2}}v_\nu + \tilde{V}_\nu v_\nu = 0$  holds in the distributional sense. Moreover  $\tilde{V}_\nu$  have the following asymptotic behaviour.

$$\tilde{V}_\nu(x) = \begin{cases} O(1/|x|), & \text{if } \frac{1}{2} < \nu < \frac{3}{2}, \nu \neq 1 \\ O(1/|x|^2), & \text{if } \nu = 1 \\ O(\log|x|/|x|), & \text{if } \nu = \frac{3}{2} \\ O(1/|x|^{4-2\nu}), & \text{if } \frac{3}{2} < \nu < 2. \end{cases} \quad (11)$$

**Remark** Since  $v_\nu(x) = O(1/|x|^{2\nu-1})$ ,  $v_\nu \in L^2(\mathbb{R})$  iff  $\nu > \frac{3}{4}$ . Thus, the massless relativistic Schrödinger operator with  $\tilde{V}_\nu$  have a zero energy eigenvalue if  $\nu > \frac{3}{4}$ , and have a zero energy resonance if  $\frac{1}{2} < \nu \leq \frac{3}{4}$ .

## References

- [1] J. Lőrinczi, I. Sasaki: Embedded Eigenvalues and Neumann-Wigner Potentials for Relativistic Schrödinger Operators, preprint, 2016
- [2] J. von Neumann, E. Wigner: Über merkwürdige diskrete Eigenwerte, *Z. Physik* **30**, 465-467, 1929