# Semiclassical analysis in the study of hyperbolic dynamical systems

Jens Wittsten

Graduate School of Human and Environmental Studies, Kyoto University

Ritsumeikan University, July 3, 2014

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 のへぐ

# Contents

- 1. Hyperbolic dynamical systems background
- 2. Semiclassical framework: How?
- 3. An explicit example studied by Faure
- 4. Results on "Spectral gap" in the semiclassical limit

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

# Dynamical systems (discrete time)

Let X be a compact Riemannian manifold and  $f: X \rightarrow X$  a smooth map.

- If f is hyperbolic then trajectories are complicated and unpredictable.
- Study instead the evolutions of densities under a linear operator induced by f.
- ► The Perron-Frobenius transfer operator M<sup>\*</sup><sub>f</sub> of f, or its relative the Ruelle transfer operator M<sub>f</sub> : φ → φ ∘ f.
- Spectral description of transfer operators can be used to obtain other specific properties of the dynamics: decay of time correlation functions, central limit theorem, mixing etc. (Ruelle, Bowen, Fried, Rugh and others).
- Spectral approach has in the last ten years been improved through the construction of functional spaces adapted to the dynamics (Blank, Gouzel, Keller, Liverani, Baladi, Tsujii and others).

# Spectral gap

What one would like: spectral gap of  $M_f : \mathcal{H} \to \mathcal{H}$ , i.e. the moduli of the largest and second largest eigenvalues (counting multiplicities) are separated by a positive distance.

 Exponential mixing is equivalent to 1 being the only eigenvalue on the unit circle.

The spectral study of the discrete spectrum of  $M_f$  (the *Ruelle resonances*) can be treated as the quantum resonances in open quantum systems.

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

Symplectic properties of the dynamics on  $T^*X$ 

ightarrow spectral properties of  $M_f$ 

 $\rightarrow$  long term behavior of the dynamics.

# Semiclassical approach

For hyperbolic dynamics, the study of transfer operators is naturally understood in the semiclassical framework. Ideas appear in

- 2007 Baladi, Tsujii.
- 2008 Baladi, Tsujii.

Formalized in

- 2006 Faure, Roy.
- 2008 Faure, Roy, Sjöstrand.

Other studies have appeared afterwards (Arnoldi, Dyatlov, Faure, Sjöstrand, Tsujii, Weich, Zworski and others).

 2011 – Faure, Semiclassical origin of the spectral gap for transfer operators of a partially expanding map, Nonlinearlity 24.

### Semiclassical approach

Find function space  $\mathcal{H}$  such that  $M_f : \mathcal{H} \to \mathcal{H}$  has a spectral gap. Note

• 
$$M_f: \varphi \mapsto \varphi \circ f$$
 is an FIO,

$$M_f\varphi(x)=\frac{1}{(2\pi)^n}\int\int e^{i(f(x)-y)\xi}\varphi(y)dyd\xi.$$

The associated symplectic map is the lift of  $f^{-1}$ ,

$$F: (f(x), \eta) \mapsto (x, {}^t f'(x)\eta).$$
(1)

▶ Usual Sobolev space H<sup>m</sup> = Op (⟨ξ⟩<sup>m</sup>)<sup>-1</sup>(L<sup>2</sup>).
 So: construct a symbol a<sub>m</sub> ∈ S<sup>m</sup> appropriately adapted to the symplectic map (1) and set H<sup>m</sup> = Op (a<sub>m</sub>)<sup>-1</sup>(L<sup>2</sup>). (Anisotropic Sobolev space.) Study M<sub>f</sub> : H<sup>m</sup> → H<sup>m</sup>.

# Proving the existence of a spectral gap I Set $A_m = Op(a_m)$ .

The commutative diagram

$$\begin{array}{cccc} L^2 & \stackrel{Q_m}{\to} & L^2 \\ \downarrow A_m^{-1} & \circlearrowleft & \downarrow A_m^{-1} \\ \mathcal{H}^m & \stackrel{M_f}{\to} & \mathcal{H}^m \end{array}$$

shows that  $M_f : \mathcal{H}^m \to \mathcal{H}^m$  is unitarily equivalent to

$$Q_m = A_m M_f A_m^{-1} : L^2 \to L^2$$

- ► Use Egorov's theorem to show that Q<sup>\*</sup><sub>m</sub>Q<sub>m</sub> is a pseudodifferential operator and calculate the principal symbol q<sub>m</sub>.
- ► Use the L<sup>2</sup> continuity theorem to get Q<sup>\*</sup><sub>m</sub>Q<sub>m</sub> = R<sub>ε</sub> + K<sub>ε</sub> where K<sub>ε</sub> is smoothing and

$$\|R_{\varepsilon}\|_{\mathcal{L}(L^{2})} \leq \limsup_{|\xi| \to \infty} \sup_{x} |q_{m}(x,\xi)| + \varepsilon, \quad \varepsilon > 0.$$

# Proving the existence of a spectral gap II

Now,

$$Q_m^* Q_m = A_m^{-1} M_f^* A_m^2 M_f A_m^{-1}$$

SO

$$q_m(x,\xi)=\frac{a_m^2\circ F(x,\xi)}{a_m^2(x,\xi)}.$$

Since  $a_m$  is strictly decreasing along trajectories of F we have  $q_m < 1$  as  $|\xi| \to \infty$ . Hence,  $Q_m^*Q_m$  has essential spectral radius < 1, and peripheral spectrum described by the compact operator  $K_{\varepsilon}$ .

Deduce the same representation for  $Q_m$  using polar decomposition and the spectral theorem. Use the unitary equivalence of  $Q_m: L^2 \to L^2$  and  $M_f: \mathcal{H}^m \to \mathcal{H}^m$  to get the same for  $M_f$ .

# The model

Let  $\mathbb{T}^2 = S^1 \times S^1$  where  $S^1 = \mathbb{R}/\mathbb{Z}$ . Define a map  $f : \mathbb{T}^2 \to \mathbb{T}^2$  by

$$f: \left(egin{array}{c} x \ s\end{array}
ight)\mapsto \left(egin{array}{c} 2x \mod 1 \ s+ au(x) \mod 1\end{array}
ight),$$

where  $E(x) := 2x \mod 1$  is uniformly expanding,

$$\min_{x} \frac{dE(x)}{dx} > 1,$$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > <

and  $\tau: S^1 \to \mathbb{R}$  is  $C^{\infty}$ . Both  $E: S^1 \to S^1$  and  $f: \mathbb{T}^2 \to \mathbb{T}^2$  are 2:1 maps.

 Assume that f is partially captive. (τ cannot be cohomologous to a constant.) Semiclassical parameter introduced by Fourier decomposition

With f as above, define the Ruelle transfer operator M<sub>f</sub> : L<sup>2</sup>(T<sup>2</sup>) → L<sup>2</sup>(T<sup>2</sup>) by

$$M_f \varphi(x, s) = \varphi(f(x, s)), \quad \varphi \in L^2(\mathbb{T}^2).$$

*M<sub>f</sub>* preserves the following decomposition in Fourier modes:

$$L^2(\mathbb{T}^2) = \bigoplus_{
u \in \mathbb{Z}} \mathcal{H}_{
u}, \quad \mathcal{H}_{
u} = \{(x,s) \mapsto \varphi(x) e^{2i\pi 
u s} : \varphi \in L^2(S^1)\}.$$

Note that for  $\psi(x,s) = \varphi(x)e^{2i\pi\nu s} \in \mathcal{H}_{\nu}$ , we have that

$$M_f\psi(x,s)=\varphi(E(x))e^{i\nu\tau(x)}e^{2i\pi\nu s}.$$

Let M<sub>ν</sub> : L<sup>2</sup>(S<sup>1</sup>) → L<sup>2</sup>(S<sup>2</sup>) denote the operator M<sub>f</sub> restricted to H<sub>ν</sub> (with identification of H<sub>ν</sub> with L<sup>2</sup>(S<sup>1</sup>)). Explicitly,

$$M_{\nu}\varphi(x) = \varphi(E(x))e^{i\nu\tau(x)}.$$

Semiclassical parameter introduced by Fourier decomposition

With f as above, define the Ruelle transfer operator M<sub>f</sub> : L<sup>2</sup>(T<sup>2</sup>) → L<sup>2</sup>(T<sup>2</sup>) by

$$M_f \varphi(x, s) = \varphi(f(x, s)), \quad \varphi \in L^2(\mathbb{T}^2).$$

► *M<sub>f</sub>* preserves the following decomposition in Fourier modes:

$$\mathcal{L}^2(\mathbb{T}^2) = \bigoplus_{\nu \in \mathbb{Z}} \mathcal{H}_{
u}, \quad \mathcal{H}_{
u} = \{(x,s) \mapsto \varphi(x) e^{2i\pi 
u s} : \varphi \in \mathcal{L}^2(\mathcal{S}^1)\}.$$

Note that for  $\psi(x,s) = \varphi(x)e^{2i\pi\nu s} \in \mathcal{H}_{
u}$ , we have that

$$M_f\psi(x,s) = \underline{\varphi(E(x))}e^{i\nu\tau(x)}e^{2i\pi\nu s}$$

Let M<sub>ν</sub> : L<sup>2</sup>(S<sup>1</sup>) → L<sup>2</sup>(S<sup>2</sup>) denote the operator M<sub>f</sub> restricted to H<sub>ν</sub> (with identification of H<sub>ν</sub> with L<sup>2</sup>(S<sup>1</sup>)). Explicitly,

$$M_{\nu}\varphi(x) = \varphi(E(x))e^{i\nu\tau(x)}.$$

For fixed  $\nu \in \mathbb{Z}$ ,

$$M_{\nu}\varphi(x) = \varphi(E(x))e^{i\nu\tau(x)}$$

is a Fourier integral operator associated to the canonical relation

$$F(y,\eta) = \{(x, E'(x)\eta) : y = E(x)\} \subset T^*(S^1)$$

When  $\nu \in \mathbb{Z}^+$  is considered as a semiclassical parameter  $h = 1/\nu$ ,

$$M_{\pm\nu}\varphi(x) = \varphi(E(x))e^{\pm i\nu\tau(x)}$$

is a semiclassical Fourier integral operator associated to the canonical relation

$$F(y,\eta) = \{(x, E'(x)\eta \pm \tau'_0(x)) : y = E(x)\} \subset T^*(S^1).$$

Since  $E: S^1 \to S^1$  is a 2:1 map, there are precisely 2 distinct elements belonging to the set in the right-hand side. Hence the map  $F: T^*S^1 \to T^*S^1$  is 2 valued.

Partial captivity essentially means that almost all trajectories escape toward infinity. For fixed  $\nu \in \mathbb{Z}$ ,

$$M_{\nu}\varphi(x) = \varphi(E(x))e^{i\nu\tau(x)}$$

is a Fourier integral operator associated to the canonical relation

$$F(y,\eta) = \{(x, E'(x)\eta) : y = E(x)\} \subset T^*(S^1)$$

When  $\nu \in \mathbb{Z}^+$  is considered as a semiclassical parameter  $h = 1/\nu$ ,

$$M_{\pm\nu}\varphi(x) = \varphi(E(x))e^{\pm i\tau(x)/h}$$

is a semiclassical Fourier integral operator associated to the canonical relation

$$F(y,\eta) = \{(x, E'(x)\eta \pm \tau'_0(x)) : y = E(x)\} \subset T^*(S^1).$$

Since  $E: S^1 \to S^1$  is a 2:1 map, there are precisely 2 distinct elements belonging to the set in the right-hand side. Hence the map  $F: T^*S^1 \to T^*S^1$  is 2 valued.

Partial captivity essentially means that almost all trajectories escape toward infinity.

#### Discrete spectrum

Recall  $\min_x E'(x) = 2$ . Theorem (Ruelle 1986, Faure 2011) Let m < 0. For any  $\nu \in \mathbb{Z}$ , the operator

 $M_{\nu}: H^m(S^1) 
ightarrow H^m(S^1)$ 

is a bounded operator which can be written

$$M_{\nu}=R_{\nu}+K_{\nu},$$

where  $K_{\nu}$  is a compact operator, and

$$\|R_
u\|_{H^m} \leq 2^m o 0, \quad m o -\infty.$$

### Spectral gap in the semiclassical limit

Let  $r(M_{\nu})$  be the spectral radius of the operator  $M_{\nu}: H^m(S^1) \to H^m(S^1)$ .

#### Theorem (Tsujii 2008, Faure 2011)

Assume that f is partially captive. Then  $r(M_{\nu})$  does not depend on m and in the semiclassical limit  $\nu \to \infty$ ,

$$r(M_{\nu}) \leq 2^{-rac{1}{2}} + o(1).$$

More precisely, for any  $\rho > 2^{-1/2}$ , there are positive constants  $m_0$  and  $\nu_0$ , such that for any  $m \ge m_0$ ,  $\nu \ge \nu_0$  and  $n \in \mathbb{N}$  we have

$$\|M_{\nu}^{n}\|_{\mathcal{L}(H_{\nu}^{m})}\leq c\rho^{n},$$

where  $H^m_{\nu}$  is the Sobolev space  $H^m(S^1)$  equipped with the norm

$$\|\varphi\|_{H^m_{\nu}}^2 = \sum_{\xi \in 2\pi\mathbb{Z}} |(1 + (\xi/\nu)^2)^m \hat{\varphi}(\xi)|^2.$$